

# NONLINEAR OPTICS (PHYC/ECE 568)

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**Homework #3, Due: Monday March 7**

## **Problem 1. Infrared Up-Conversion**

Estimate the efficiency of the upconversion of 10  $\mu\text{m}$  infrared radiation using sum frequency generation. The pump laser has  $\lambda=532$  nm, with a power of 10 W. Use 2 cm long Proustite (use  $d_{\text{eff}}=d_{22}$ ) crystal under perfect phase-matching and optimum focusing where  $L=2Z_0$  ( $Z_0$  is the Rayleigh range of the focused laser beam).

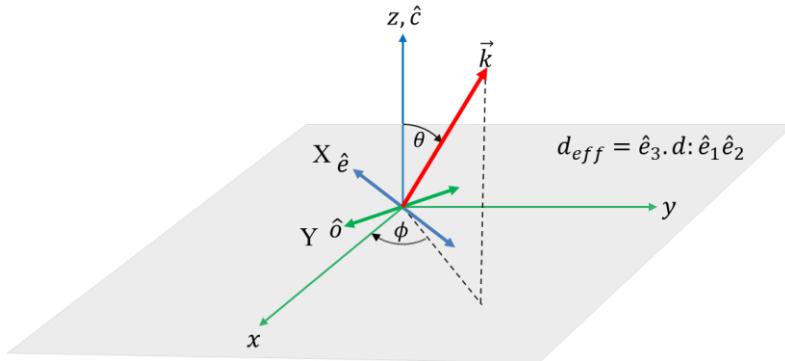
## **Problem 2.**

(a) Show that  $d_{\text{eff}}$ , as defined by  $P_3 = 4d_{\text{eff}} E_1 E_2$ , is related to the  $d$  tensor via:

$$d_{\text{eff}} = \hat{e}_3 \bullet d : \hat{e}_1 \hat{e}_2$$

where  $\hat{e}_j$  ( $j=1,2,3$ ) is the unit vector associated with  $E_1, E_2$ , and  $P_3$ .

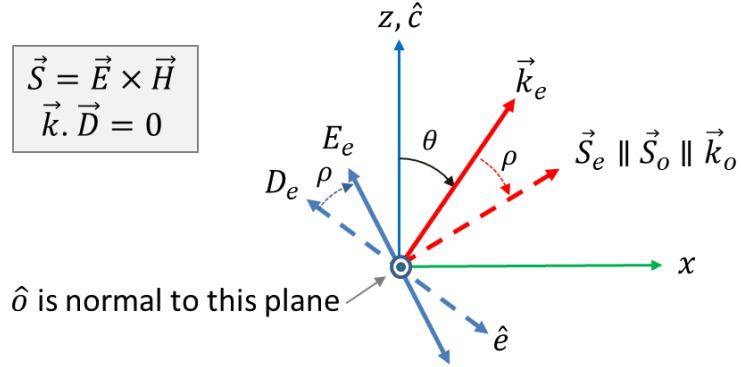
(b) For a given geometry,  $d_{\text{eff}}$  is usually calculated in terms of  $d_{il}$ 's and the angles  $\phi$  and  $\theta$  as described in the figure below. Here  $x, y$  and  $z$  (or 1,2 and 3) are the crystal axis and X, Y, and Z (laboratory frame) are optical propagation axis (e.g.  $\mathbf{k}_2 \parallel \mathbf{k}_1 \parallel \mathbf{Z}$ ). Note: Y is in xy-plane (thus normal to z- or optics axis) and X is on zZ plane.



- (i) Derive expressions for  $d_{\text{eff}}$  for a class 3m crystal (e.g. LiNbO<sub>3</sub>) where  $\hat{e}_1 = \hat{e}_2 = Y$  (ordinary), and  $\hat{e}_3 = X$  (extra-ordinary). (This is known as type-I condition).
- (ii) Repeat the above calculation for type-II condition where  $\hat{e}_1 = Y$  (ordinary) ,  $\hat{e}_2 = X$  and  $\hat{e}_3 = X$  (extra-ordinary).
- (iii) Find a geometry (i.e.  $\theta$  and  $\phi$ ) that accesses the largest  $d_{il}$  element in LiNbO<sub>3</sub>. (see data provided here).
- (iv) Find the phase matching angle for SHG generation at  $\lambda$  (fundamental)=1.15  $\mu\text{m}$  for the part (i) and (ii). Discuss the phase matching situation for case (iii)

**Problem 3. Poynting Vector Walk-off:**

We know, from linear optics, that the e- and o-rays in a birefringent crystal walk-off from each other (i.e. double-refraction) resulting from the fact that  $k_e$  and  $k_o$  are not parallel. Known as Poynting vector walk-off, this is essentially the angle ( $\rho$ ) between E and D vectors for the e-ray where  $D=\epsilon: E$ .



In the harmonic generation applications, such as SHG, this imposes a serious restriction on the useful length of the nonlinear crystal.

a. Assuming a uniaxial crystal, calculate the walk-off angle  $\rho$  between e- and o-ray Poynting vectors. Show that

$$\tan \rho = -\frac{1}{n_e(\theta)} \frac{dn_e(\theta)}{d\theta}$$

b. Show that for type-I phase matching SHG (o+o $\rightarrow$ e)

$$\rho \approx \tan \rho = \frac{n_o^2(\omega)}{2} \left[ \frac{1}{\tilde{n}_e^2(2\omega)} - \frac{1}{n_o^2(2\omega)} \right] \sin(2\theta_m)$$

## LiNbO<sub>3</sub> Properties

$$\begin{bmatrix} 0 & 0 & 0 & 0 & d_{31} & -d_{22} \\ -d_{22} & d_{22} & 0 & d_{31} & 0 & 0 \\ d_{31} & d_{31} & d_{33} & 0 & 0 & 0 \end{bmatrix} \text{ class 3m crystal}$$

## Nonlinear Optical Coefficients of LiNbO<sub>3</sub> at 1.06 μ m

$d_{22} / 1 \text{ d}_{36}^{\text{KDP}} 1$	6.5
$d_{31} / 1 \text{ d}_{36}^{\text{KDP}} 1$	-12.3
$d_{33} / 1 \text{ d}_{36}^{\text{KDP}} 1$	-86

$$d_{36} \text{ (KDP).} = 0.4 \text{ pm/V}$$

## Refractive Indices at 20°C

Wavelength, μm	n <sub>o</sub>	n <sub>e</sub>
0.43584	2.39276	2.29278
0.54608	2.31657	2.22816
0.63282	2.28647	2.20240
1.1523	2.2273	2.1515
3.3913	2.1451	2.0822

# NLO - PS # 2 Solution

Problem ① →

(Conversion (quantum) efficiency defines as :

$$\begin{cases} n_1 = 1 - \left| \frac{A_1(L)}{A_1(0)} \right|^2 \quad \text{or} \\ n_2 = \left| \frac{A_3(L)}{A_1(0)} \right|^2 = n_1 \frac{\omega_3}{\omega_1} \end{cases}$$

$$\text{Boyd (2.4.7)} \rightarrow n_1 = 1 - \cos^2(KL) = \sin^2(KL)$$

$$K = \frac{8\pi d}{c} \sqrt{\frac{\omega_1 \omega_3}{n_1 n_3}} |A_2| \rightarrow$$

- in order to have  $K$  in SI, we should set  $\begin{cases} 4\pi \epsilon_0 \rightarrow 1 \\ d \rightarrow \epsilon_0 d \end{cases}$

$$\rightarrow K = \frac{2d}{c} \sqrt{\frac{\omega_1 \omega_3}{n_1 n_3}} |A_2| \Rightarrow K^2 = \frac{2d^2 \omega_1 \omega_3}{c^3 \epsilon_0 n_1 n_2 n_3} I_2$$

$$\rightarrow \text{Also, we know : } I_2 = 2n_2 \epsilon_0 c |A_2|^2$$

$$I_2 = \frac{P_2}{\text{Area}} = \frac{2P_2}{\pi \omega_0^2},$$

$$\text{for optimum focusing : } L = 2\tau_0 = \frac{2\pi n_2 \omega_0}{\lambda_2} = \frac{n_2 \omega_2 \omega_0}{c}$$

$$\rightarrow \text{From above relations : } I_2 = \frac{2P_2 n_2 \omega_2}{\pi c L}$$

$$\Rightarrow K^2 L^2 = \frac{4d^2 \omega_1 \omega_2 \omega_3}{\pi c^4 \epsilon_0 n_1 n_3} P_2 L = \frac{32 d^2 P_2 L}{\pi \epsilon_0 c \lambda_1 \lambda_2 \lambda_3 n_1 n_3}; \lambda_3 = \frac{\lambda_1 \lambda_2}{\lambda_1 + \lambda_2}$$

$$\Rightarrow K^2 L^2 = \frac{32 d^2 P_2 L (\lambda_1 + \lambda_2)}{\pi \epsilon_0 c (\lambda_1 \lambda_2)^2 n_1 n_3}$$

→ based on given value

$$n_1 = \frac{K^2 L^2}{c^2} \approx 3.197 \%$$

(for  $d = 36 \text{ mm}$ )

problem (2)  $\rightarrow$

(a)

$$\vec{P}_3 = 4d : \vec{E}_1 \vec{E}_2$$

$\hat{P}_3 \hat{e}_3 = 4d : \hat{e}_1 \hat{e}_2 \vec{E}_1 \vec{E}_2 \rightarrow$  Apply  $(\hat{e}_3 \cdot)$  on both sides:

$$P_3 (\hat{e}_3 \cdot \hat{e}_3) = 4 \hat{e}_3 \cdot d : \hat{e}_1 \hat{e}_2 \vec{E}_1 \vec{E}_2$$

$$\rightarrow P_3 = 4d_{\text{eff}} \vec{E}_1 \vec{E}_2 ; \text{ where } d_{\text{eff}} = \hat{e}_3 \cdot d : \hat{e}_1 \hat{e}_2 \quad \checkmark$$

b-i

$$\hat{e}_1 = \hat{e}_2 = X \quad \& \quad \hat{e}_3 = Y$$

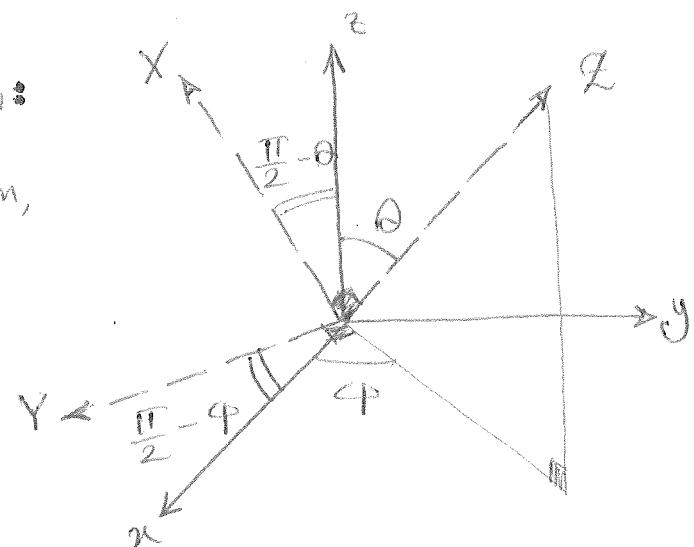
$$\int X = -\cos \theta \cdot \cos \varphi \hat{x} + \cos \theta \sin \varphi \hat{y} + \sin \theta \hat{z}$$

$$\int Y = \sin \varphi \hat{x} - \cos \varphi \hat{y}$$

\*  $d : \hat{e}_1 \hat{e}_2$  is evaluated as below:

- after applying  $\hat{e}_3 \cdot ( )$  operation,

there only  $\hat{x}$  &  $\hat{y}$  components of  
 $d : \hat{e}_1 \hat{e}_2$  remain!



$$\begin{bmatrix} d_{11} & d_{12} & d_{13} & d_{14} & d_{15} & d_{16} \\ d_{21} & d_{22} & d_{23} & d_{24} & d_{25} & d_{26} \\ d_{31} & d_{32} & d_{33} & d_{34} & d_{35} & d_{36} \end{bmatrix} \begin{bmatrix} \cos^2 \theta & \cos^2 \varphi \\ \cos^2 \theta & \sin^2 \varphi \\ \sin^2 \theta \\ -2 \sin \theta \cos \theta \sin \varphi \\ -2 \sin \theta \cos \theta \cos \varphi \\ 2 \cos^2 \theta \sin \varphi \cos \varphi \end{bmatrix}$$

$$\begin{aligned}
 \rightarrow d_{\text{eff}} = & \cos^2 \theta \cos^2 \varphi (d_{11} \sin \varphi - d_{21} \cos \varphi) + \\
 & \cos^2 \theta \sin^2 \varphi (d_{12} \sin \varphi - d_{22} \cos \varphi) + \\
 & \sin^2 \theta \quad (d_{13} \sin \varphi - d_{23} \cos \varphi) + \\
 & - 2i \sin \theta \sin \varphi (d_{14} \sin \varphi - d_{24} \cos \varphi) + \\
 & - 2i \sin \theta \cos \varphi (d_{15} \sin \varphi - d_{25} \cos \varphi) + \\
 & \cos^2 \theta \sin 2\varphi (d_{16} \sin \varphi - d_{26} \cos \varphi) .
 \end{aligned}$$

\* for an 3m crystal :

$$d = \begin{vmatrix} 0 & 0 & 0 & 0 & d_{31} & -d_{22} \\ -d_{22} & d_{22} & 0 & d_{31} & 0 & 0 \\ d_{31} & d_{31} & d_{33} & 0 & 0 & 0 \end{vmatrix}$$

$$\begin{aligned}
 \rightarrow d_{\text{eff}} = & [\cos^2 \theta \cos^2 \varphi \cos \varphi - \cos^2 \theta \sin^2 \varphi \cos \varphi - \\
 & \cos^2 \theta \sin 2\varphi \sin \varphi] d_{22} + \\
 & + [\sin 2\theta \sin \varphi \cos \varphi - \sin 2\theta \cos \varphi \sin \varphi] d_{31} \\
 = & \cos^2 \theta [\cos \varphi (\cos 2\varphi) - \sin 2\varphi \sin \varphi] d_{22} \\
 = & \cos^2 \theta \cos 3\varphi d_{22}
 \end{aligned}$$

- for the case of negative birefringent crystal,

$$\hat{e}_1 = \hat{e}_2 = \hat{x}$$

$$\hat{e}_3 = \hat{x}$$

(Haw question)

$\mathbf{J} \cdot \hat{\mathbf{e}}_1 \hat{\mathbf{e}}_2 =$

$$= \begin{pmatrix} d_{11} & d_{12} & \dots & d_{16} \\ d_{21} & d_{22} & \dots & d_{26} \\ d_{31} & d_{32} & \dots & d_{36} \end{pmatrix} \begin{pmatrix} 2i^2 \varphi \\ Cn^2 \varphi \\ 0 \\ 0 \\ 0 \\ -2i \theta Cn \varphi \end{pmatrix}$$

$$= \begin{pmatrix} d_{11} 2i^2 \varphi + d_{12} Cn^2 \varphi - d_{16} 2i \theta \\ d_{21} 2i^2 \varphi + d_{22} Cn^2 \varphi - d_{26} 2i \theta \varphi \\ d_{31} 2i^2 \varphi + d_{32} Cn^2 \varphi - d_{36} 2i \theta \varphi \end{pmatrix} \rightsquigarrow \text{apply } (\hat{\mathbf{e}}_3 \cdot) \rightarrow$$

$$\begin{aligned} d_{\text{eff}} = & (d_{11} 2i^2 \varphi + d_{12} Cn^2 \varphi - d_{16} 2i \theta \varphi) (-Cn \theta Cn \varphi) + \\ & (d_{21} 2i^2 \varphi + d_{22} Cn^2 \varphi - d_{26} 2i \theta \varphi) (-Cn \theta 2i \varphi) + \\ & (d_{31} 2i^2 \varphi + d_{32} Cn^2 \varphi - d_{36} 2i \theta \varphi) (2i \theta). \end{aligned}$$

For 3m crystal:

$$\begin{aligned} d_{\text{eff}} = & -d_{22} [2i^2 \varphi Cn \theta Cn \varphi + (2i^2 \varphi - Cn^2 \varphi) (-Cn \theta 2i \varphi)] + \\ & + d_{31} [2i^2 \varphi + Cn^2 \varphi] 2i \theta. \end{aligned}$$

$$= -d_{22} [Cn \theta (Cn 2i \varphi + 2i^2 \varphi Cn \varphi)] + d_{31} 2i \theta$$

$$\rightarrow d_{\text{eff}} = -d_{22} Cn \theta 2i^3 \varphi + d_{31} 2i \theta \quad \boxed{\checkmark}$$

b-ii

we can consider 2 cases:  $\begin{cases} \text{Case-1 : } \hat{e}_1 = \hat{x}, \hat{e}_2 = \hat{y}, \hat{e}_3 = \hat{x} \\ \text{Case-2 : } \hat{e}_1 = \hat{x}, \hat{e}_2 = \hat{y}, \hat{e}_3 = \hat{y} \end{cases}$

$\Rightarrow \hat{e}_1, \hat{e}_2 =$

$$\begin{bmatrix} d_{11} & \dots & d_{16} \\ d_{21} & \dots & d_{26} \\ d_{31} & \dots & d_{36} \end{bmatrix} \begin{bmatrix} -Cn\theta \text{ Cn}2\varphi \text{ Zn}2\varphi \\ Cn\theta \text{ Zn}2\varphi \text{ Cn}2\varphi \\ 0 \\ 0 \\ 0 \\ Cn\varphi \text{ Cn}\theta - Zn\varphi \text{ Cn}\theta \end{bmatrix} = \begin{bmatrix} \frac{1}{2}Cn\theta Zn2\varphi (d_{12} - d_{11}) + Cn\theta Cn2\varphi d_{16} \\ \frac{1}{2}Cn\theta Zn2\varphi (d_{22} - d_{21}) + Cn\theta Cn2\varphi d_{26} \\ \frac{1}{2}Cn\theta Zn2\varphi (d_{32} - d_{31}) + Cn\theta Cn2\varphi d_{36} \end{bmatrix}$$

\* applying  $(\hat{e}_3 = \hat{x})$  :

$\Rightarrow \hat{e}_3 = \hat{x} : \text{ Case 1}$

$$\begin{aligned} d_{\text{eff}} &= \frac{1}{2} Cn\theta Zn2\varphi \left[ -Cn\theta Cn\varphi (d_{12} - d_{11}) - Cn\theta Zn\varphi (d_{22} - d_{21}) + Zn\theta (d_{32} - d_{31}) \right. \\ &\quad \left. + Cn\theta Cn2\varphi \left[ -Cn\theta Cn\varphi d_{16} - Cn\theta Zn\varphi d_{26} + Zn\theta d_{36} \right] \right] \end{aligned}$$

$\Rightarrow \hat{e}_3 = \hat{y} : \text{ Case 2}$

$$\begin{aligned} d_{\text{eff}} &= \frac{1}{2} Cn\theta Zn2\varphi \left[ Zn\varphi (d_{12} - d_{11}) - Cn\varphi (d_{22} - d_{21}) \right] \\ &\quad + Cn\theta Cn2\varphi \left[ Zn\varphi d_{16} - Cn\varphi d_{26} \right] \end{aligned}$$

\* for a 3m crystal :

$$\begin{aligned} \text{Case-1 : } d_{\text{eff}} &= \frac{1}{2} Cn\theta Zn2\varphi Cn\theta Zn\varphi (-2d_{22}) + \\ &\quad - Cn\theta Cn2\varphi Cn\theta Cn\varphi (-d_{22}) \\ &= d_{22} Cn^2\theta Cn3\varphi \quad \rightarrow \text{H-W question} \end{aligned}$$

$$\begin{aligned}
 \text{Case-2 : } d_{\text{eff}} &= -\frac{1}{2} \text{ Crd } \sin \theta \text{ Crp } (2d_{22}) + \\
 &\quad \text{ Crd } \sin 2\theta \text{ Crp } (-d_{22}) \\
 &= -d_{22} \text{ Crd } \sin 3\theta
 \end{aligned}$$

⑥ - (iii)

- to access  $d_{33}$  (largest in  $\text{LiNbO}_3$ ), one needs :

$$\begin{aligned}
 \hat{e}_1 \parallel \hat{e}_2 \parallel \hat{e}_3 \parallel X \parallel z \\
 \Rightarrow \text{ if } \theta = \frac{\pi}{2} \Rightarrow d_{\text{eff}} = d_{33}
 \end{aligned}$$

- we shall see that this configuration is not phase-matched using the crystal birefringent. But, employing a growth technique known as "Periodic Poling" we can enhance phase matching under above geometry!

⑥ - (iv)

$$\lambda_1 = 1.15 \mu\text{m} , \quad \lambda_2 = \frac{\lambda_1}{2} = 0.570$$

we know for,

$$\lambda = 0.632 \rightarrow n_o = 2.2864 , \quad n_e = 2.2024$$

$$\lambda = 0.546 \rightarrow n_o = 2.3165 , \quad n_e = 2.2281$$

- Using linear interpolation of indices at these two wavelengths, we'll obtain  $n_o$  &  $n_e$  for  $\lambda_2 = 0.570$  :

$$n_o(\lambda_2) = 2.306$$

$$n_e(\lambda_2) = 2.219$$

Type-I PM :

$$n_e(\theta_m, \lambda_2) = n_o(\lambda_1) \rightarrow$$

$$2n^2 \theta_m = \frac{n_o^{-2}(\lambda_1) - n_o^{-2}(\lambda_2)}{n_e^{-2}(\lambda_2) - n_o^{-2}(\lambda_2)} = 0.906 \rightarrow \theta_m = 72^\circ$$

Type-II PM :

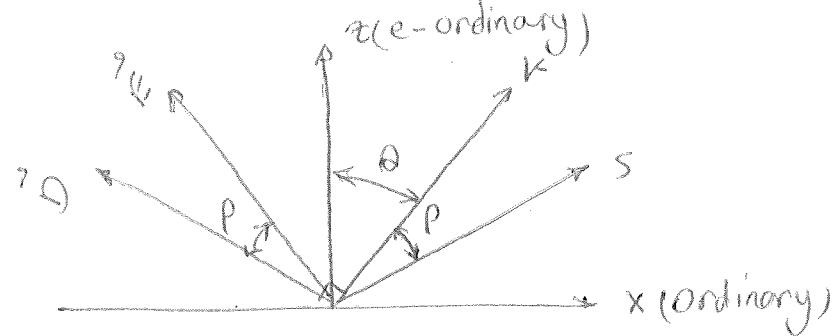
$$n_o(\lambda_1) + n_e(\theta_m, \lambda_1) = \frac{\lambda_1}{\lambda_2} n_e(\theta_m, \lambda_2)$$

- Upon inserting the numbers, one finds (using either graphical or math softwares) that there is no solution to this. That is, this condition is not type-II phase matchable.

Accessing  $d_{33}$  required that  $\theta = 90^\circ$ . This condition is not phase matchable for this wavelength. Quasi Phase Matching (QPM) is ideal as has been used in periodically poled lithium niobate (PPLN).

### problem ③

The poynting vector walk-off for a uniaxial crystal can be treated in a two-dimensional analysis.



Consider the extra-ordinary ray. (i.e.  $\theta \neq 0$ ), for this cone,  $\vec{k} \perp \vec{D}$  &  $\vec{S} \perp \vec{E}$ , the walk-off angle  $\rho$  is also the angle between  $\vec{E}$  &  $\vec{D}$ !

$$\rightarrow \vec{D} \cdot \vec{E} = |\vec{D}| \cdot |\vec{E}| \cos \rho \rightarrow \cos \rho = \frac{\vec{D} \cdot \vec{E}}{|\vec{D}| \cdot |\vec{E}|}$$

$$\rightarrow \cos \rho = \frac{D_x E_x + D_z E_z}{(D_x^2 + D_z^2)^{1/2} (E_x^2 + E_z^2)^{1/2}}$$

$$\text{but } D_x = -D \cos \theta$$

$$D_z = D \sin \theta$$

$$\text{also } D_x = \epsilon_x E_x = n_0^2 E_x$$

$$D_z = \epsilon_z E_z = n_e^2 E_z$$

$$\rightarrow E_x = -\frac{D \cos \theta}{n_0^2} \quad \text{and} \quad E_z = \frac{D \sin \theta}{n_e^2}$$

$$\rightarrow \cos \rho = \frac{D^2 \left[ \frac{\cos^2 \theta}{n_0^2} + \frac{\sin^2 \theta}{n_e^2} \right]}{D^2 \left[ \frac{\cos^2 \theta}{n_0^4} + \frac{\sin^2 \theta}{n_e^4} \right]^{1/2}} \rightarrow \cos^2 \rho = \frac{\left( \frac{\cos^2 \theta}{n_0^2} + \frac{\sin^2 \theta}{n_e^2} \right)^2}{\frac{\cos^2 \theta}{n_0^4} + \frac{\sin^2 \theta}{n_e^4}}$$

$$\Rightarrow \tan^2 \rho = \frac{1}{Cn^2 \rho} - 1$$

$$= \left( \frac{Cn^2 \theta}{n_b^2} + \frac{2i^2 \theta}{n_e^2} \right)^{-2} \left[ \frac{Cn^2 \theta}{n_b^4} + \frac{2i^2 \theta}{n_e^4} - \frac{Cn^4 \theta}{n_b^4} - \frac{2i^4 \theta}{n_e^4} - \frac{22i^2 Cn^2 \theta}{n_b^2 n_e^2} \right]$$

$$\Rightarrow \tan \rho = \frac{Cn \theta \sin \left( \frac{1}{n_b^2} - \frac{1}{n_e^2} \right)}{\left( \frac{Cn^2 \theta}{n_b^2} + \frac{2i^2 \theta}{n_e^2} \right)}$$

$$= \frac{1}{2} \left[ \frac{Cn^2 \theta}{n_b^2} + \frac{2i^2 \theta}{n_e^2} \right]^{-1} \left( \frac{1}{n_b^2} - \frac{1}{n_e^2} \right) 2i^2 \theta \quad \checkmark$$

(b)

for Type - I SHG, where  $0 + 0 \rightarrow e$

$$\frac{1}{n_b^2(\omega)} = \frac{Cn^2 \theta_m}{n_b^2(2\omega)} + \frac{2i^2 \theta_m}{n_e^2(2\omega)}$$

thus,

$$\tan \rho = \frac{1}{2} n_b^2(\omega) \times \left[ \frac{1}{n_b^2(2\omega)} - \frac{1}{n_e^2(2\omega)} \right] \sin 2\theta_m \quad \checkmark$$

Solution for second part of 3(a) in HW22\_3 :

$$\frac{1}{n_e(\theta)^2} = \frac{\sin(\theta)^2}{\tilde{n}_e^2} + \frac{\cos(\theta)^2}{n_o^2}$$

Take derivative with respect to  $\theta$ :

$$\frac{-2 \frac{dn_e(\theta)}{d\theta}}{n_e(\theta)^3} = 2 \sin(\theta) \cos(\theta) \left( \frac{1}{\tilde{n}_e^2} - \frac{1}{n_o^2} \right) = \sin(2\theta) \left( \frac{1}{\tilde{n}_e^2} - \frac{1}{n_o^2} \right)$$

Therefore

$$-\frac{1}{n_e(\theta)} \frac{dn_e(\theta)}{d\theta} = \frac{1}{2} n_e(\theta)^2 \sin(2\theta) = \tan(\rho)$$