

# Laser Physics I (PHYC/ECE 464)

Homework #4, Due Monday, Sept. 26, Fall 2022

1

Consider a linear combination of two equal amplitude  $\text{TEM}_{m,p}$  modes given by:

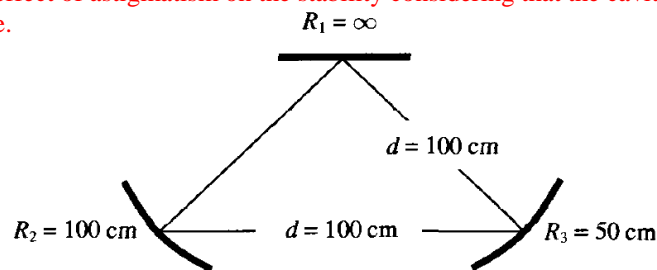
$$\mathbf{E} = E_0 \{ (\text{TEM}_{1,0})\mathbf{a}_y \pm j(\text{TEM}_{0,1})\mathbf{a}_x \}$$

- (a) Sketch the “dot” pattern or equal intensity contours for each component (i.e.,  $\mathbf{a}_x$  or  $\mathbf{a}_y$ ). Indicate the direction of the electric field.
- (b) Sketch the pattern for the linear combination.
- (c) Label the positions where the intensity is a maximum and a minimum. (This is sometimes referred to as the “donut mode” or  $\text{TEM}_{0,1}^*$ ).

2

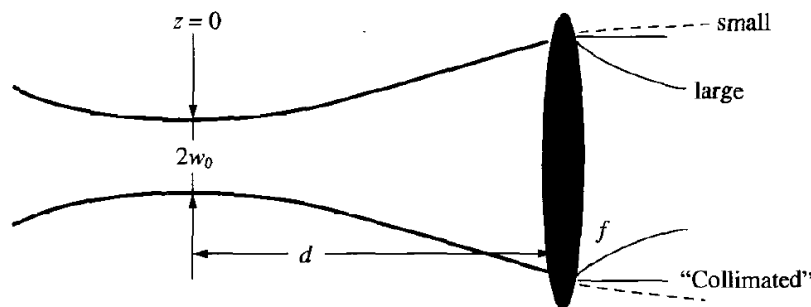
Is the cavity shown below stable? Demonstrate the logic of your answer by (a) constructing a unit cell starting at the flat mirror, (b) finding the  $ABCD$  matrix for that cell, and (c) applying the stability criteria. (d) What are the circumstances under which the quantity  $[AD - BC]$  can be different from 1? Why is  $AD - BC$  always equal to 1 for a cavity? (Ignore astigmatism)

**Bonus:** Show the effect of astigmatism on the stability considering that the cavity axis below forms an equilateral triangle.



3

A focused Gaussian beam reaches its minimum spot size  $w_0$  at  $z = 0$  where  $R = \infty$  and then propagates to a thin lens of focal length  $f$  located a distance  $d$  from  $z = 0$ . If  $w_0$  is large, then the beam exiting the lens will be focused. If it is too small, then the lens merely reduces the far field spreading angle. Find the critical value of  $w_0$  such that the output beam is “collimated”; that is,  $R(z = d^+) = \infty$  also.



4. Using Eq. (5.2.8), write  $z_0^2$  in terms of the g-parameters (i.e.  $g_1$  and  $g_2$ ). From this, derive the stability condition of the cavity in terms of  $g_1$  and  $g_2$ . Compare this with the geometric optics results obtained earlier.

①

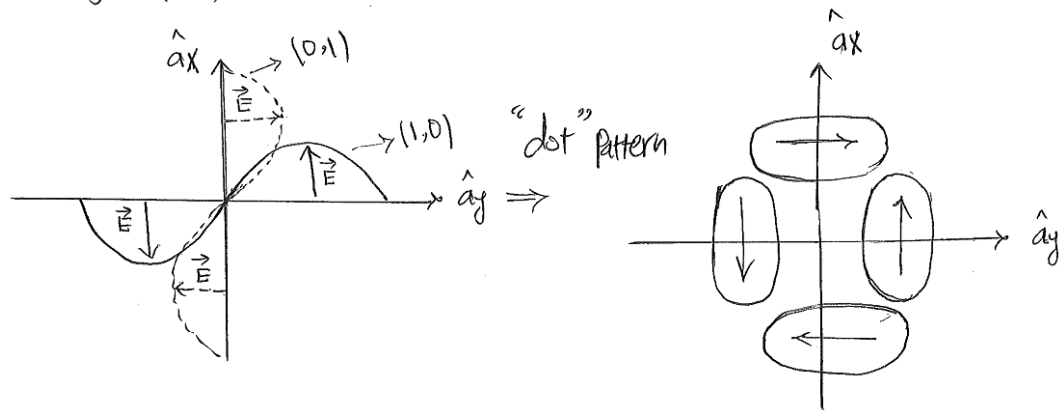
Problem (3.6) Verdeyen 3<sup>rd</sup> ed.

$$\vec{E} = E_0 \left( (\text{TEM}_{1,0}) \hat{a}_y \pm j (\text{TEM}_{0,1}) \hat{a}_x \right)$$

$$(a) E(x,y,z) = E_0 \rho H_m \left( \frac{\sqrt{2}x}{w(z)} \right) H_p \left( \frac{\sqrt{2}y}{w(z)} \right) \frac{w_0}{w(z)} e^{\frac{-(x^2+y^2)}{w(z)^2}} e^{-j(\text{phase})}$$

↳ Equation 3.5.1 of the text.

and we know that  $H_0(u) = 1$  &  $H_1(u) = 2u$  and according to the figure (3.5) of the text,

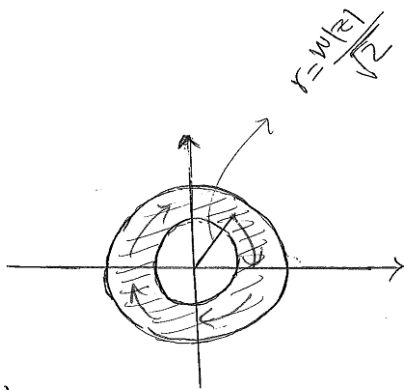


$$(b) |E|^2 = |E_0|^2 \left( (\text{TEM}_{1,0})^2 + (\text{TEM}_{0,1})^2 \right)$$

$$\Rightarrow |E|^2 = |E_0|^2 \left( \frac{2}{2} \left( \frac{\sqrt{2}x}{w(z)} \right)^2 x^2 \left( \frac{w_0}{w(z)} \right)^2 e^{\frac{-2x^2}{w(z)^2}} + x^2 \left( \frac{\sqrt{2}y}{w(z)} \right)^2 \left( \frac{w_0}{w(z)} \right)^2 e^{\frac{-2y^2}{w(z)^2}} \right)$$

$$= |E_0|^2 4 \left( \frac{\sqrt{2}}{w(z)} \right)^2 \left( \frac{w_0}{w(z)} \right)^2 (x^2 + y^2) e^{\frac{-2r^2}{w(z)^2}}$$

$$= \alpha r^2 e^{\frac{-2r^2}{w(z)^2}}$$



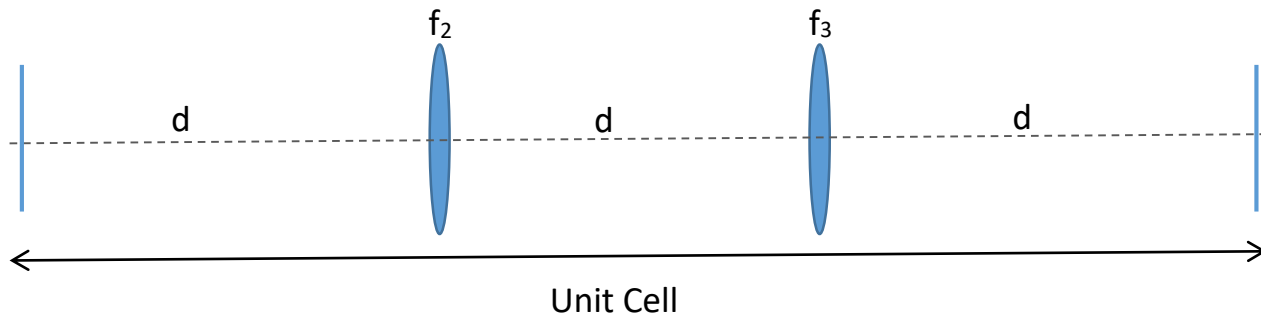
$$I \sim |E|^2 \sim r^2 e^{-\frac{2r^2}{w(z)^2}}$$

$$\frac{dI}{dr} \sim \frac{d}{dr} \left( \alpha r^2 e^{-\frac{2r^2}{w(z)^2}} \right) = \alpha \left( 2r e^{-\frac{2r^2}{w(z)^2}} - \frac{4r}{w(z)^2} r^2 e^{-\frac{2r^2}{w(z)^2}} \right) = 0$$

$$\Rightarrow 1 - \frac{2r^2}{w(z)^2} = 0 \Rightarrow r = \frac{w(z)}{\sqrt{2}} \quad (\text{Maximum occurs at this radius}).$$

$$\left( \frac{d^2 I}{dr^2} \right)_{r = \frac{w(z)}{\sqrt{2}}} < 0$$

Verdeyen 3.18: Ignoring the astigmatism introduced by the oblique incidence on curved mirrors



The equivalent transmission system for the ring cavity is

$$T = \begin{pmatrix} 1 & d \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & d \\ -1/f_3 & 1 - d/f_3 \end{pmatrix} \begin{pmatrix} 1 & d \\ -1/f_2 & 1 - d/f_2 \end{pmatrix} = \begin{pmatrix} 1 - \frac{2d}{f_2} - \frac{d}{f_3} + \frac{d^2}{f_2 f_3} & \text{something} \\ \text{something} & 1 - \frac{2d}{f_3} - \frac{d}{f_2} + \frac{d^2}{f_2 f_3} \end{pmatrix}$$

If  $d=100$  cm,  $f_2=50$ cm,  $f_3=25$  cm, then  $A=+1$  and  $D=-1$ , and  $S = \frac{A+D+2}{4} = \frac{1}{2}$  ( $>0$  and  $<1$ ); the cavity is therefore stable.

If the entrance and exit planes have different indices of refraction, the  $AD-BC=n_2/n_1 \neq 1$ .

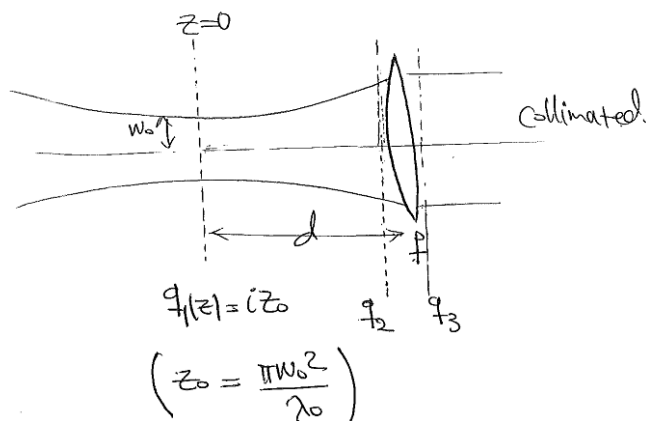
In a cavity (roundtrip unit cell), the start and stop planes are always the same, thus  $AD-BC=1$  always.

**BONUS PART:** With astigmatism, the angle of incidence on each mirror is  $\theta=60/2$  degrees. Then,  $f$ 's above will need to change to  $f/\cos(\theta)$  and  $f^*\cos(\theta)$  for sagittal and tangential components respectively.

This leads to  $S=0.103$  (for sagittal, stable) and  $S=1.137$  (tangential, unstable). The cavity then as a whole is not stable.

③

Problem (3.21) (Verdeyen 3<sup>rd</sup> ed.)



The radius of the beam at  $z=d$  is,

$$R(z) = z \left( 1 + \left( \frac{z_0}{z} \right)^2 \right) \Rightarrow R(z=d) = d \left( 1 + \left( \frac{z_0}{d} \right)^2 \right) = \frac{z_0^2 + d^2}{d}$$

The transformation matrix for a thin lens  $M = \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix}$

$$\frac{1}{q_3} = \frac{C + D(1/q_2)}{A + B(1/q_2)} = \frac{-1/f + 1/q_2}{1 + 0} = -\frac{1}{f} + \frac{1}{q_2}$$

$$= \frac{1}{f} + \frac{1}{R_2} - j \frac{\lambda_0}{\pi w_2^2}$$

on the other hand,

$$\frac{1}{q_3} = \frac{1}{R_3} - j \frac{\lambda_0}{\pi w_3^2} = \frac{1}{f} + \frac{1}{R_2} - j \frac{\lambda_0}{\pi w_2^2}$$

$$\Rightarrow \frac{1}{R_3} = \frac{1}{f} + \frac{1}{R_2} \Rightarrow R_3 = \frac{R_2 f}{R_2 - f} \xrightarrow{R_3 \rightarrow \infty} R_2 = f$$

$$R_2 = \frac{d+z^2}{d} = f \Rightarrow z^2 = fd - d^2$$

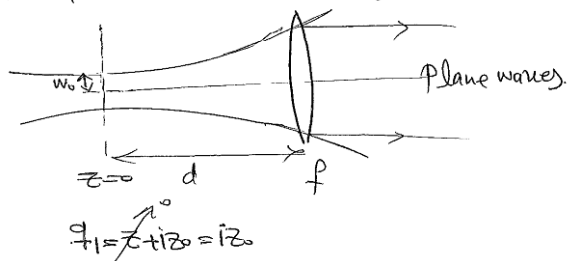
$$z_0 = \frac{\pi w_0^2}{\lambda_0} ; \quad w_0^4 = \frac{\lambda_0^2}{\pi^2} (fd - d^2)$$

$$\Rightarrow w_0 = \sqrt{\frac{\lambda_0}{\pi}} (d|f-d|)^{1/4}$$

So the beam waist must be this amount to have a plane wave after the lens. ( $R_3 = \infty$ )

Alternative way:

calculate the matrix from the beam waist  $z=0$  to the lens and then use the ABCD law to find the  $q$  parameter after the lens.



$$T = \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix} \begin{pmatrix} 1 & d \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & d \\ -\frac{1}{f} & 1 - \frac{d}{f} \end{pmatrix}$$

Now, use  $\frac{1}{q_2} = \frac{C+D(1/q_1)}{A+B(1/q_1)}$

$$\begin{aligned}\Rightarrow \frac{1}{q_2} &= \frac{-1/f + (1-d/f) \frac{1}{iz}}{1 + d/iz} = \frac{-iz_0/f + (1-d/f)}{iz_0 + d} \times \frac{-iz_0 + d}{iz_0 + d} \\ &= \frac{d(1-d/f) - z_0^2/f - iz_0 d/f - iz_0(1-d/f)}{d^2 + z_0^2} \\ &= \frac{d(1-d/f) - z_0^2/f - i(z_0 d/f + z_0 - z_0 d/f)}{d^2 + z_0^2}\end{aligned}$$

on the other hand,  $\frac{1}{q_2} = \frac{1}{R_2(z)} - i \frac{\lambda_0}{\pi W_2(z)}$

Equating  $\frac{1}{q_2}$ 's we have,  
for the real part

$$\Rightarrow \frac{1}{R_2(z)} = \frac{d(1-d/f) - z_0^2/f}{d^2 + z_0^2} \Rightarrow R_2(z) = \frac{d^2 + z_0^2}{d(1-d/f) - z_0^2/f}$$

denominator must go to zero to have  $R_2 \rightarrow \infty$ .

$$d(1-d/f) - \frac{z_0^2}{f} = 0 \Rightarrow z_0^2 = d(f-d) \quad (z_0 = \frac{\pi W_0^2}{\lambda_0})$$

$$\Rightarrow W_0^4 = \frac{\lambda_0^2}{\pi^2} d(f-d) \Rightarrow W_0 = \left( \frac{\lambda_0^2}{\pi^2} d(f-d) \right)^{1/4}$$

④

$$z_0^2 = \frac{d(R_1-d)(R_2-d)(R_1+R_2+d)}{(R_1+R_2-2d)^2} ; g_{1,2} = 1 - \frac{d}{R_{1,2}} = \frac{R_{1,2}-d}{R_{1,2}}$$

$$\Rightarrow z_0^2 = \frac{d(R_1 g_1)(R_2 g_2)(R_1+R_2-d)}{(g_1 R_1 + g_2 R_2)^2} =$$

$$= \frac{d^2 R_1 R_2 g_1 g_2 (1-g_1 g_2)}{(g_1 R_1 + g_2 R_2)^2 (1-g_1)(1-g_2)}$$

$$= \frac{d^4 \left(\frac{1}{1-g_1}\right) \left(\frac{1}{1-g_2}\right) g_1 g_2 (1-g_1 g_2)}{(1-g_1)(1-g_2) \left(\frac{g_1 d}{1-g_1} + \frac{g_2 d}{1-g_2}\right)^2} = \frac{d^4 g_1 g_2 (1-g_1 g_2)}{(1-g_1)^2 (1-g_2)^2 \left[d(g_1+g_2-2g_1 g_2)\right]^2}$$

$$\Rightarrow z_0^2 = \frac{d^2 g_1 g_2 (1-g_1 g_2)}{(g_1+g_2-2g_1 g_2)^2}$$

Stability Cond.  $z_0^2 \geq 0 \Rightarrow g_1 g_2 (1-g_1 g_2) \geq 0$

$$\Rightarrow \begin{cases} g_1 g_2 \geq 0 \\ 1-g_1 g_2 \geq 0 \end{cases} \Rightarrow 0 \leq g_1 g_2 < 1$$