

Formula Sheet    PHYC/ECE 463    Advanced Optics I    UNM/Fall 2007

harmonic plane wave:  $E = \text{Re}\{E_0 \exp(i\phi)\}$  with  $\phi \equiv \omega t - k \cdot r + \varphi$ ,  $k = \frac{n\omega}{c} = \frac{2\pi n}{\lambda_0}$

Poynting vector  $S = \frac{1}{\mu_0} E \times B$     Irradiance:  $I = \langle S \rangle = \frac{\epsilon_0 n c}{2} |E_0|^2$

$n_i \sin(\theta_i) = n_t \sin(\theta_t)$  Snell's Law

$$\rho_\pi = \frac{n_t \cos(\theta_i) - n_i \cos(\theta_t)}{n_t \cos(\theta_i) + n_i \cos(\theta_t)} = \frac{\tan(\theta_i - \theta_t)}{\tan(\theta_i + \theta_t)} \quad \tau_\pi = \frac{2n_i \cos(\theta_i)}{n_t \cos(\theta_i) + n_i \cos(\theta_t)}$$

$$\rho_\sigma = -\frac{n_i \cos(\theta_i) - n_t \cos(\theta_t)}{n_i \cos(\theta_i) + n_t \cos(\theta_t)} = -\frac{\sin(\theta_i - \theta_t)}{\sin(\theta_i + \theta_t)} \quad \tau_\sigma = \frac{2n_i \cos(\theta_i)}{n_i \cos(\theta_i) + n_t \cos(\theta_t)}$$

$$1 - \rho^2 = \tau\tau' \quad R = |\rho|^2 \quad T = |\tau|^2 \frac{n_t \cos(\theta_t)}{n_i \cos(\theta_i)} \quad R+T=1$$

If  $n$  is complex then  $n \rightarrow \tilde{n} = n - i\kappa \equiv \sqrt{\epsilon' / \epsilon_0} \equiv \sqrt{1 + \chi}$  in above expressions

Absorption coefficient ( $K$  or  $\alpha$ ) and skin depth ( $\delta$ ):  $\alpha \equiv \frac{2}{\delta} = \frac{4\pi\kappa}{\lambda_0}$

Classical Electron Oscillator Model:

$$\chi = \frac{\omega_p^2}{\omega_0^2 - \omega^2 + i\omega/\tau} \quad \text{where } \omega_p = \sqrt{\frac{Nq^2}{m_0\epsilon_0}} \text{ (plasma frequency)}$$

Drude model for metals:  $\omega_0 \rightarrow 0$

Group Velocity  $v_g = \frac{d\omega}{dk}$

Light pressure (on perfecting absorbing surface)  $P = \frac{I}{c}$

Prism with apex angle  $\alpha$  at minimum deviation angle  $\theta_D$ :

$$n = \frac{\sin\left(\frac{\alpha + \theta_D}{2}\right)}{\sin\left(\frac{\alpha}{2}\right)}$$

Numerical Aperture (NA) of an optical fiber:

$f\#$  ( $f$ -number) =  $f/D$

$$n_0 \sin(\theta_{\max}) = \sqrt{n_f^2 - n_c^2}$$

Lens-makers' formula:      Gaussian imaging formula (thin lens)      (refractive sphere)

$$\frac{1}{f} = (n-1) \left( \frac{1}{R_1} - \frac{1}{R_2} \right) \qquad \frac{1}{f} = \frac{1}{S} + \frac{1}{S'} \qquad \frac{n'-n}{R} = \frac{n}{S} + \frac{n'}{S'}$$

Paraxial Ray Tracing Matrices:

Propagation $\begin{pmatrix} 1 & 0 \\ d/n & 1 \end{pmatrix}$ (length d and index n)	Refractive surface $\begin{pmatrix} 1 & -P \\ 0 & 1 \end{pmatrix}$ where $P=(n'-n)/R$
Mirror with radius of curvature R $\begin{pmatrix} 1 & -2/R \\ 0 & 1 \end{pmatrix}$	Thin lens of focal length f $\begin{pmatrix} 1 & -1/f \\ 0 & 1 \end{pmatrix}$
Thick lens (Thickness $D_l$ , index $n_l$ ) $\begin{pmatrix} 1 - \frac{P'D_l}{n_l} & -P - P' + \frac{PP'D_l}{n_l} \\ \frac{D_l}{n_l} & 1 - \frac{PD_l}{n_l} \end{pmatrix}$ $P=(n_l-n)/R, \quad P'=(n'-n_l)/R'$	Separated Doublet $\begin{pmatrix} 1 - \frac{d}{f_2 n_b} & -\frac{1}{f_1} - \frac{1}{f_2} + \frac{d}{f_2 f_1 n_b} \\ \frac{d}{n_b} & 1 - \frac{d}{f_1 n_b} \end{pmatrix}$

between conjugate planes $\begin{pmatrix} m_\alpha \frac{n'}{n} & \tilde{M}_{12} \\ 0 & m_x \end{pmatrix}$	between principal planes $\begin{pmatrix} 1 & M_{12} \\ 0 & 1 \end{pmatrix}$
for telescopic Systems $\begin{pmatrix} m_\alpha \frac{n'}{n} & 0 \\ M_{21} & m_x \end{pmatrix}$	position of principal planes $D = \begin{pmatrix} n \\ M_{12} \end{pmatrix} (1 - M_{11})$ $D' = \begin{pmatrix} n' \\ M_{12} \end{pmatrix} (1 - M_{22})$

Contact Doublet Achromatization:

$$f_1 V_1 + f_2 V_2 = 0 \quad \text{where } V = \frac{n_d - 1}{n_f - n_c} \text{ is the Abbe number}$$

Physical constants:

speed of light	$c \sim 2.998 \times 10^8$	$\text{m s}^{-1}$
electronic charge	$q \sim 1.602 \times 10^{-19}$	C
permittivity of vacuum	$\epsilon_0 \sim 8.854 \times 10^{-12}$	F/m
electronic mass	$m_0 \sim 9.1094 \times 10^{-31}$	Kg
Planck constant	$h \sim 6.626 \times 10^{-34}$	J.s

Near distance of a normal eye:  $d_0=250 \text{ mm}$

## Two-beam Interference

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\delta) \quad \text{where } \delta = \phi_1 - \phi_2$$

## Multiple-Beam Interference:

$$\text{Wavefront splitting (N slits, waves having equal amplitudes): } I(\delta) = I(0) \left( \frac{\sin(N\delta/2)}{N \sin(\delta/2)} \right)^2$$

$$\delta = 2\pi \frac{a}{\lambda} (\sin \theta' - \sin \theta) \quad a = \text{slit separation}, \quad \delta = 2m\pi \quad (\text{constructive interference})$$

$$\text{Resolving power } \mathcal{R}: \frac{\lambda}{\Delta\lambda} = mN$$

$$\text{Amplitude Splitting: } \delta = \frac{4\pi}{\lambda_0} n_2 d \cos \theta_2 - \Delta\gamma \quad \Delta\gamma = \text{roundtrip reflection phase (if any)}$$

$$\text{Fabry-Perot: (for } R_1=R_2=R) \quad T = \frac{T_{\max}}{1 + F \sin^2(\delta/2)} \quad \text{where } T_{\max} = \left(1 - \frac{A}{1-R}\right)^2 \quad \text{and } A = \text{loss (mirror)}$$

$$F = \frac{4R}{(1-R)^2} \quad (\text{contrast}), \quad \mathcal{R} = \frac{\lambda}{\Delta\lambda} = m\mathcal{F} \quad \text{where } \mathcal{F} = \frac{\pi\sqrt{F}}{2} \approx \frac{\pi}{1-R} \quad \text{is the Finesse}$$

$$\text{as a (laser) cavity: } \mathcal{F} = \frac{\Delta\nu_{FSR}}{\Delta\nu_{1/2}} \quad \Delta\nu_{FSR} = c/2nd$$

## Matrix Formalism (thin film coatings):

$$E_j = \begin{pmatrix} E_{lj} \\ E_{rj} \end{pmatrix} \quad \text{fields on the left side of layer } j, \quad E'_j = \begin{pmatrix} E'_{lj} \\ E'_{rj} \end{pmatrix} \quad \text{fields on the right side of layer } j$$

$$\text{interface transition: } E'_i = H_{ij} E_j \quad \text{where } H_{ij} = \frac{1}{\tau_{ij}} \begin{pmatrix} 1 & \rho_{ij} \\ \rho_{ij} & 1 \end{pmatrix}$$

$$\text{propagation: } E_j = L_j E'_j \quad \text{where } L_j = \begin{pmatrix} e^{-i\beta_j} & 0 \\ 0 & e^{+i\beta_j} \end{pmatrix} \quad \text{with } \beta_j = \frac{2\pi}{\lambda_0} n_j d_j \cos \theta_j \quad (\text{note } n_j \text{ may be}$$

complex if loss or gain is present).

For N-2 layers (sandwiched between medium 1 and N):

$$S = H_{12} L_2 H_{23} L_3 \dots L_{N-1} H_{N-1,N} = \begin{pmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{pmatrix}, \quad \tau = \frac{1}{S_{22}}, \quad \rho = \frac{S_{12}}{S_{22}}$$

## Waveguides:

$$\text{Assuming: } E(r,t) = \mathcal{E}(x,y) e^{i\omega t - i\beta z} \quad \text{then} \quad \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \mathcal{E}(x,y) + (n^2 k_0^2 - \beta^2) \mathcal{E}(x,y) = 0^{**}$$

$$\text{Slab waveguide solutions: core: } \mathcal{E}(x) = \begin{cases} A \sin(hx) \\ A \cos(hx) \end{cases} \quad \text{cladding: } \mathcal{E}(x) = B \begin{cases} e^{-p(x-d/2)} \\ e^{+p(x+d/2)} \end{cases} \quad \text{where}$$

$$h = \sqrt{n_2^2 k_0^2 - \beta^2} \quad p = \sqrt{\beta^2 - n_1^2 k_0^2}$$

\*  $h$  and  $p$  are obtained from conditions that tangential components ( $E_y$  and  $H_z \propto \partial E_y / \partial x$  for TE,  $H_y$  and  $E_z \propto (1/n^2) \partial H_y / \partial x$  for TM) are continuous at the interfaces.

\*\* For fibers the above Helmholtz equation\*\* is solved in cylindrical coordinates.  $E(r, \phi, t) = \mathcal{E}(r) e^{-i\ell\phi} e^{i\omega t - i\beta z}$

$$\mathcal{E}(r) \propto \begin{cases} J_\ell(k_r r) & r < a(\text{core}) \\ K_\ell(\gamma r) & r > a(\text{cladding}) \end{cases} \quad \sqrt{(k_r a)^2 + (\gamma a)^2} = 2\pi \frac{a}{\lambda_0} (n_2^2 - n_1^2) = V \quad (\text{V-number})$$

$$\text{Number of modes } M \approx \frac{4}{\pi^2} V^2 \quad (\text{step index fibers})$$