

**Nonlinear Optics (PHYC/ECE 568)**  
**Spring 2007**

*Midterm Exam I, Open Text-Book, Time: 1.5 hour*

*NAME* .....  
*last* ..... *first*

Grade
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*Please staple and return this page with your exam.*

**1. SHG Bandwidth (a HW problem):**

**a.** Calculate the bandwidth  $\Delta\omega$  associated with a phase-matched SHG process in terms of the group velocities  $v_g(\omega_1)$  and  $v_g(2\omega_1)$ . *Bandwidth is defined as the frequency spread around  $\omega_1$ , for which  $\Delta k(\omega_1 \pm \Delta\omega/2)L = 2\pi$ .* with  $L$  denoting the length of the nonlinear crystal. *Hint: Use the first-order term in the Taylor series expansion of  $\Delta k(\omega)$ .* (15 pts.)

**b.** Discuss how your results in (a) explains the limitation on the SHG-efficiency when ultrashort laser pulses are used.

(5 pts.)

Note group velocity  $v_g = (dk/d\omega)^{-1}$

**2. SHG Acceptance Angle:**

**a.** In a similar manner as in problem 1, calculate the acceptance angle ( $\Delta\theta$ ) for a type-I phase matched SHG process. *Acceptance angle is defined as angular spread around  $\theta_m$ , for which  $\Delta k(\theta_m \pm \Delta\theta/2)L = 2\pi$ .* Assume negative birefringence in a uniaxial crystal where  $\Delta k(\theta) = (2\omega/c)(n_e^{2\omega}(\theta) - n_o^\omega)$ . (20 pts.)

**b.** Assuming small birefringence ( $n_e \approx n_o$ ), show that

$$\Delta\theta \cong \frac{\lambda_\omega}{L(n_o^{2\omega} - n_e^{2\omega}) \sin(2\theta_m)} \quad (5 \text{ pts.})$$

**c.** What limitation does this impose on the useful length of the crystal for an incident focused Gaussian beam of waist  $w_0$  (recall: divergence angle is  $\lambda/\pi n w_0$ ). (10 pts.)

**d.** Briefly discuss the implications of the above relationship for various angles (e.g. for  $\theta_m = 90^\circ$ ). (5 pts.)

**3.  $d_{\text{eff}}$ :**

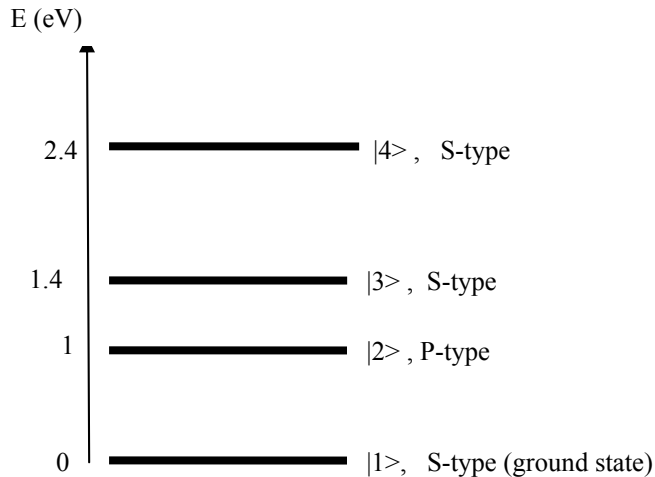
Consider  $\text{AgGaSe}_2$  crystal in a SHG process. The orientation of the input field at  $\omega$  is given by  $E(\omega) = A_0(2\hat{x} + 1\hat{y} + 3\hat{z})$ , where x, y, and z are the principle axes of the crystal. What is  $d_{\text{eff}}$ ? (10 pts.)

Note: From table 1.5.3 and Fig. 1.5.2 (in Boyd) you get  $d_{ij}$  (pm/V) tensor in  $\text{AgGaSe}_2$  to be:

$$\begin{pmatrix} 0 & 0 & 0 & 81 & 0 & 0 \\ 0 & 0 & 0 & 0 & 81 & 0 \\ 0 & 0 & 0 & 0 & 0 & -81 \end{pmatrix}$$

4. A fictional molecule has the following energy levels. Draw the spectrum ( for  $0 < \hbar\omega < 3 \text{ eV}$ ) for the (a) linear absorption coefficient  $\alpha$ , (b) two-photon absorption coefficient  $\beta$ , (c)  $|\chi^{(2)}(2\omega; \omega, \omega)|$  and (d)  $|\chi^{(3)}(3\omega; \omega, \omega, \omega)|$ . (30 pts.)

Be quantitative in your x-axis. Assume a finite broadening in your drawings. Point out the resonances (diagrammatically) on your graph for each case and show the relative strengths if obvious. Note: no calculations needed for this problem.



1. See H.W. Solution

2. SHG B.W

$$n_o > n_e$$

$$\Delta k(\theta) = \frac{2\omega}{c} (n_e^{2\omega}(\theta) - n_o^{\omega})$$

Taylor Series Expansion

$$\Delta k(\theta) = \Delta k(\theta_m) + \left. \frac{\partial k}{\partial \theta} \right|_{\theta=\theta_m} (\theta - \theta_m) + \frac{1}{2} \left. \frac{\partial^2 k}{\partial \theta^2} \right|_{\theta=\theta_m} (\theta - \theta_m)^2 + \dots$$

$$\Delta k(\theta_m + \frac{\Delta\theta}{2}) L = \left. \frac{\partial k}{\partial \theta} \right|_{\theta_m} \frac{\Delta\theta}{2} L = 2\pi$$

$$\Delta\theta = \frac{4\pi}{\left( \left. \frac{\partial k}{\partial \theta} \right|_{\theta_m} \right) L}$$

$$\frac{\partial k}{\partial \theta} = \frac{2\omega}{c} \frac{\partial n_e^{2\omega}(\theta)}{\partial \theta}$$

$$(n_e^{2\omega}(\theta))^{-2} = (n_e^{2\omega})^{-2} \sin^2 \theta + (n_o^{2\omega})^{-2} \cos^2 \theta$$

$$n_e^{2\omega}(\theta) = \left[ (n_e^{2\omega})^{-2} \sin^2 \theta + (n_o^{2\omega})^{-2} \cos^2 \theta \right]^{-\frac{1}{2}}$$

$$\left. \frac{\partial n_e^{2\omega}}{\partial \theta} \right|_{\theta=\theta_m} = -\frac{1}{2} \left[ \frac{2 \sin \theta_m \cos \theta_m}{(n_e^{2\omega})^2} + \frac{-2 \sin \theta_m \cos \theta_m}{(n_o^{2\omega})^2} \right] \left[ \right]^{-\frac{3}{2}}$$

$$= -\frac{\Delta \epsilon \omega^2 \Delta m}{2} \left[ \frac{1}{(n_e^{\omega})^2} - \frac{1}{(n_o^{\omega})^2} \right] \cdot \left[ \quad \right]^{-3/2}$$

$(n_o^{\omega})^3$

$$\frac{\partial n_e(\omega)}{\partial \Delta} = \frac{-\Delta \epsilon \omega^2 \Delta m}{2} \cdot (n_o^{\omega})^3 \left( \frac{1}{(n_e^{\omega})^2} - \frac{1}{(n_o^{\omega})^2} \right)$$

(b)

$$n_e \approx n_o = n \quad \Delta n$$

$$\frac{\partial n_e(\omega)}{\partial \Delta} \approx -\frac{\Delta \epsilon \omega^2 \Delta m}{2} \cdot \frac{n^3}{n^4} (n_o^2 - n_e^2)^{2\omega} = \text{~~... \Delta n~~}$$

$$\approx -\frac{\Delta \epsilon \omega^2 \Delta m}{2} \cdot \frac{n^3}{n^4} 2n \Delta n$$

$$\approx \Delta \epsilon \omega^2 \Delta m \Delta n$$

$$\Delta n \equiv (n_o - n_e)_{2\omega}$$

$$\frac{\partial k}{\partial \Delta} \Big|_{\omega} = \frac{2\omega}{c} \Delta \epsilon \omega^2 \Delta m \Delta n = \frac{4\pi \Delta n \Delta \epsilon \Delta m}{\lambda_{\omega}}$$

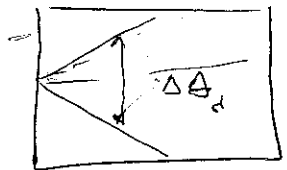
$$DQ = \frac{\Delta m}{\Delta n \Delta \epsilon \Delta m \cdot L}$$

(c)

$$\Delta Q_{div} = \frac{\lambda}{\pi n W_0}$$

we must have

$$\Delta \theta_{div} < 2 \Delta \theta$$



Note:  $\Delta \theta_{div}$  is half-angle  
 $\Delta \theta$  is full angle  
by definition

$$\frac{\lambda_w}{\pi n W_0} < \frac{2 \lambda_w}{\Delta n \sin^2 \theta_m \cdot L}$$

$$L < \frac{2 \pi n W_0}{\Delta n \sin^2 \theta_m} = l_d$$

(Note this is <sup>almost</sup> the same as  $l_a$  from Rayleigh vector walk-off). (HW#3)

$$l_a = \frac{\sqrt{\pi n W_0}}{\Delta n \sin^2 \theta_m}$$

(d)

$$\text{at } \theta_m = \frac{\pi}{2}, \quad l_d \rightarrow \infty \downarrow$$

This is because  $\frac{\partial k}{\partial \theta} \rightarrow 0$  and we must take higher order <sup>terms</sup> in the expansion into account. ( $\frac{\partial^2 k}{\partial \theta^2}, \dots$ ).

$$\text{This leads to } \Delta \theta = \sqrt{\frac{\partial w}{2L (n_o^{2w} - n_e^{1w})}}$$

3

$$E(t) = A_0 (2\hat{x} + 1\hat{y} + 3\hat{z})$$

$$= \underbrace{A_0 \sqrt{14}}_{E_0} \left( \frac{2}{\sqrt{14}}\hat{x} + \frac{1}{\sqrt{14}}\hat{y} + \frac{3}{\sqrt{14}}\hat{z} \right)$$

$\hat{e}$  unit vector.

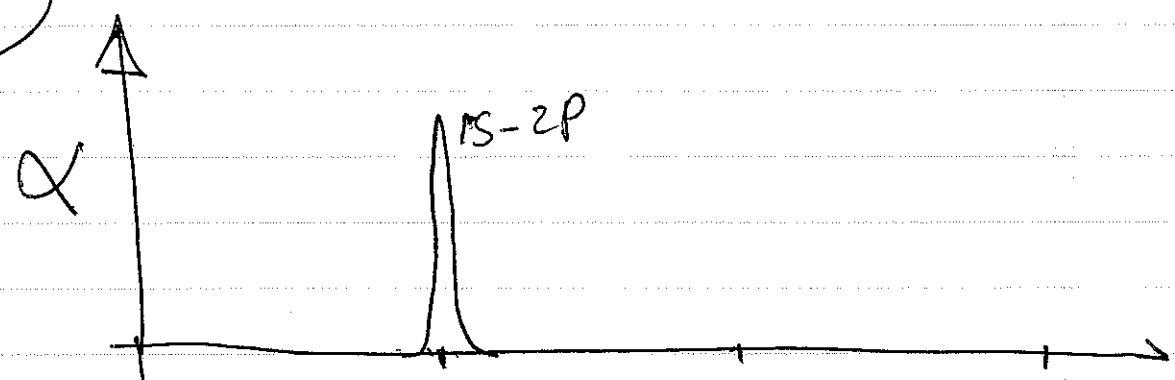
$$P = 2 \epsilon_0 \epsilon_b |E_0|^2$$

$$= 2 \cdot \begin{bmatrix} 0 & 0 & 0 & 81 & 0 & 0 \\ 0 & 0 & 0 & 0 & -81 & 0 \\ 0 & 0 & 0 & 0 & 0 & -81 \end{bmatrix} \begin{bmatrix} E_x \\ E_y \\ E_z \\ 2E_y E_z \\ 2E_x E_z \\ 2E_y E_x \end{bmatrix}$$

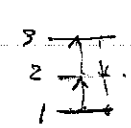
$$\begin{bmatrix} \frac{4}{14} \\ \frac{1}{14} \\ \frac{9}{14} \\ \frac{6}{14} \\ \frac{12}{14} \\ \frac{4}{14} \end{bmatrix} |E_0|^2$$

$$\epsilon_b = 81 \left[ \frac{6}{14} - \frac{12}{14} - \frac{4}{14} \right] = -\frac{810}{14} \text{ perm/V}$$

(4)

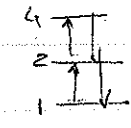


~~β~~  
 $g_m \{ \chi^{(2)} \}$



1S-3S

0.7



1S-4S

1.2

1.2 eV is close to 1 eV.  
 (1ω) resonance.

$\chi^{(1)}$

$\chi^{(2)} \equiv 0$  (Symmetry)

$\chi^{(3)}$   
 $\chi^{(3\omega)}$

(1-2) 3ω  
 (1-3) 2ω  
 (1-2) 1ω  
 (1-4) 2ω

0.33

0.7

1.2

1

2

3

eV